

1.
(i) Let $u^{(∞)}(x) = Ax + B$

$$\begin{cases} u^{(∞)}(0) = B = 0 \\ u^{(∞)}(L) = AL = 0 \end{cases}$$

$$u^{(∞)}(L) = 0 \quad \# \quad +5$$

(ii) Assume $u(x,t) = X(x)T(t)$

$$\frac{X''}{X} = \frac{1}{\alpha^2} \frac{T'}{T} = -k^2$$

$$\Rightarrow \begin{cases} X(x) = C \cos kx + D \sin kx \\ T(t) = E e^{-\alpha^2 k^2 t} \end{cases} \quad \Leftarrow +5$$

using boundary condition

$$\Rightarrow u(x,t) = \sum_{n=1}^{\infty} A_n \sin k_n x e^{-\alpha^2 k_n^2 t}, \quad k_n = \frac{n\pi}{L}$$

$$A_n = \frac{2}{L} \int_0^L f(x) \sin \frac{n\pi}{L} x dx \quad \Leftarrow +10$$

↑
I.C.

$$A_n = \begin{cases} \frac{160}{n^2\pi^2}, & n=1, 5, 9, \dots \\ -\frac{160}{n^2\pi^2}, & n=3, 7, 11, \dots \\ 0, & \text{others} \end{cases} \quad \Leftarrow +15$$

(iii) 速度:

改變後的 $k' = \frac{n\pi}{2L} \Rightarrow$ 較慢

或寫一邊過熱, \Rightarrow 較慢

對一個 +2

全對 +5

new steady-state:

$$\begin{cases} u_x^{(∞)}(0) = A = 0 \\ u_x^{(∞)}(L) = B = 0 \end{cases} \Rightarrow u_x^{(∞)} \text{ 不變}$$

(iv)

兩邊取熱 \Rightarrow steady-state 為原熱能均勻分佈

$$u^{(\infty)}(x) = \frac{20 \times L}{2} \times \frac{1}{L} = 10$$

① 寫出不為零的 constant

② 寫為 10, (+5)

2 (v)

Let $u^{(p)}(x, 0) = \delta(x)$

$$\Rightarrow u^{(p)}(x, t) = \frac{1}{2\alpha\sqrt{4\pi t}} e^{-\frac{x^2}{4\alpha^2 t}}$$

← 有看到這個就 (+3)

then, $u(x, t) = f(x) * u^{(p)}(x, t)$ (+5) (through x)

$$= \int_{-\infty}^{\infty} f(\tau) u^{(p)}(x-\tau, t) d\tau$$

$$u(x, t) = \int_{-\infty}^{\infty} f(\tau) u^{(p)}(-\tau, t) d\tau$$

完整寫出 convolution (+1)

$$= \frac{-u_0 \alpha \sqrt{t}}{a \sqrt{\pi}} \left(e^{-\frac{b^2}{4\alpha^2 t}} - e^{-\frac{a^2}{4\alpha^2 t}} \right)$$

(+10)

3. (i) 有寫 $c = f\lambda + 2$ (ii) 沒化簡 $\sum_{n=1}^{\infty} + 7$ ~~列對算~~

單位 ~~6m+3~~ ~~嘗試計算+1~~

寫

寫 $W = c\pi \sqrt{\left(\frac{L}{3}\right)^2 + \left(\frac{m}{6}\right)^2 + \left(\frac{n}{3}\right)^2}$ ~~Ans對~~

化成 346 \square ~~沒~~

$w = 2\pi f$ ~~沒~~

$f = \frac{w}{2\pi}$ rad \square

60 Hz \rightarrow 3分

$\sum C_n \sin \frac{n\pi x}{3} \sin \frac{m\pi y}{6} \sin \frac{p\pi z}{3}$

\downarrow
不扣, 因為可被 T 吸收

x, y, z 順序錯不扣

~~分~~
 $\sin \alpha x \sin \beta y \sin \gamma z$
(general form only) + 5

x, y, z 寫成 x, x, x ~~最~~

$\frac{c^2 U''}{U} = \frac{T''}{T} = -k^2$ ~~分~~

A D F = 0 \square + 2

(mnp) (1,1,1) + 2

16

標錯 \square 或沒標 \square ~~標~~

2-2 \square 滿分

N_{sym} 不扣分

$(2, 2, 2) + 5$ \square $\frac{3}{2} - 1$

$u_p(\omega) = 0 + \omega \rightarrow \omega \rightarrow \infty = 0$

$u_p(\omega) = 0 = \omega \rightarrow \infty = 0 + 8$

\square + 10

\square + 8 \square + 8

全 0 \square sloppy 扣分

\square + 3

4. (i) We can expand a well-behaved function to an orthogonal and complete basis. $\sin(n\frac{\pi}{L}x)$ is a complete basis, we can write $g(x,t) = \sum_{n=1}^{\infty} G_n(t) \sin(n\frac{\pi}{L}x)$

(ii)
$$y(x,t) = \sum_{n=1}^{\infty} b_n(t) \sin(n\frac{\pi}{L}x)$$

$$\left\{ \begin{aligned} c^2 \frac{\partial^2 y}{\partial x^2} &= \frac{d^2 y}{dt^2} + a \frac{dy}{dt} - g(x,t) \end{aligned} \right.$$

$$\Rightarrow -c^2 \sum_{n=1}^{\infty} \left(\frac{n\pi}{L}\right)^2 b_n(t) \sin(n\frac{\pi}{L}x)$$

$$= \sum_{n=1}^{\infty} b_n''(t) \sin(n\frac{\pi}{L}x) + \sum_{n=1}^{\infty} a b_n'(t) \sin(n\frac{\pi}{L}x)$$

$$- \sum_{n=1}^{\infty} G_n(t) \sin(n\frac{\pi}{L}x)$$

\Rightarrow Compare the coefficient of $\sin(n\frac{\pi}{L}x)$

$$c^2 \left(\frac{n\pi}{L}\right)^2 b_n(t) + b_n''(t) + a b_n'(t) = G_n(t)$$

(iii)
$$\left\{ \begin{aligned} y(x,0) = 0 &\Rightarrow y(x,0) = \sum_{n=1}^{\infty} b_n(0) \sin(n\frac{\pi}{L}x) = 0 \\ y_t(x,0) = 0 &\Rightarrow y_t(x,0) = \sum_{n=1}^{\infty} b_n'(0) \sin(n\frac{\pi}{L}x) = 0 \end{aligned} \right.$$

Imply $b_n(0) = b_n'(0) = 0$

From (ii)'s solution

$$b_n''(t) + a b_n'(t) + \omega_n^2 b_n(t) = G_n(t)$$

where $\omega_n = \frac{cn\pi}{L}$

For this kind of solution may have e^{-at} term,

but $t \gg \frac{1}{a}$, it will disappear

有富叶开式
+ 2 ~ t5

尝试解

in general solution

Therefore, we just need to solve steady-state solution.

Try $b_n(t) = A e^{i\omega t}$ and $G_n(t) = g_0 \cos \omega t = g_0 \operatorname{Re}[e^{i\omega t}]$

$$-A\omega^2 e^{i\omega t} + aA i\omega e^{i\omega t} + \omega_n^2 A e^{i\omega t} = g_0 e^{i\omega t}$$

$$\Rightarrow A(-\omega^2 + i a \omega + \omega_n^2) = g_0$$

$$\Rightarrow A = \frac{g_0}{(\omega_n^2 - \omega^2) + i a \omega}$$

We can express A in exponential notation $|A| e^{-i\phi}$

$$|A| = \frac{g_0}{\sqrt{(\omega_n^2 - \omega^2)^2 + a^2 \omega^2}}$$

$$\tan \phi = \frac{a\omega}{\omega_n^2 - \omega^2}$$

Then we have

$$b_n(t) = \operatorname{Re}[|A| e^{i\phi} e^{i\omega t}]$$

$$= \frac{g_0}{\sqrt{(\omega_n^2 - \omega^2)^2 + a^2 \omega^2}} \cos(\omega t - \phi)$$

where $\tan \phi = \frac{a\omega}{\omega_n^2 - \omega^2}$



Laplace equation
 Laplace
 $t > \frac{1}{a}$

後面針壓
 可也
 $\frac{1}{a}$

(iv) Because it has "ay_c" term, waves will decay.

The shape will decrease exponentially and it can't be expressed by two waves reflection in the walls, that is

$$F(x+ct) + G(x-ct)$$

No → 3

$$\frac{P}{\omega \lambda + (\omega - \omega)} = A$$

No + (8) → 5

Yes → 0

Yes + (8) → 1 ~ 2

$$\frac{P}{\omega \lambda + (\omega - \omega)} = |A|$$

$$\frac{P}{\omega \lambda + (\omega - \omega)} = P_{inc}$$

$$\frac{P}{\omega \lambda + (\omega - \omega)}$$

$$P_{inc} = |A| |A| = (A)^2$$

$$\cos(\omega t - \phi)$$

$$\frac{\omega}{\omega - \omega}$$

$$P_{inc} = P_{ref} = P_{tr}$$

$$\cos(\omega t - \phi) = \sin(\omega t + \phi) = \cos(\phi - \omega t)$$