

Final (Modern Physics) (6 problems on both sides of this paper)

06/11/2018 Provided by Masahito Oh-e

Instructions:

- You are not allowed to open any textbook, copies of my lecture notes and ppt files, but allowed to see **your own notebook or memo**. You can also use **a simple calculator**.
- Do not use the Internet. If anyone is found who is cheating on the exam, she/he will be immediately failed in this course.
- Solve the problems below. Describe the ways of thinking in English: **only final solutions are not accepted**. Make clear how you reach each solution.
- When 30 minutes pass after the test starts, if you think you have completed the test, you can leave the room by submitting the answer sheets, **except the 10 minutes before the exam finishes**.
- Several physical constants are listed in the end of the sheet.
- If answers are decimal numbers, calculate them to one places of decimals at least.

Problem 1. (25 points)

Consider the particle described by a wavefunction: $\psi = e^{ix} + 2ie^{3ix}$ ($-\pi \leq x \leq \pi$)

- Normalize the wavefunction.
- If you precisely measure the momentum of the state expressed by the wavefunction ψ , what values can you obtain and what probability can you get, respectively?
- Calculate the probability of finding the particle as a function of the position.
- Draw the graph of the probability density.
- Calculate the expectation value of the position.

$$(a) \psi^* \psi = (e^{-ix} - 2ie^{-3ix})(e^{ix} + 2ie^{3ix}) = 5 - 2ie^{-2ix} + 2ie^{2ix}$$

$$\int_{-\pi}^{\pi} \psi^* \psi dx = \int_{-\pi}^{\pi} (5 - 2ie^{-2ix} + 2ie^{2ix}) dx = 10\pi$$

Therefore, the normalized function is $\psi = \frac{1}{\sqrt{10\pi}} (e^{ix} + 2ie^{3ix})$

$$\begin{aligned} & (e^{-ix} - 2ie^{-3ix})(e^{ix} + 2ie^{3ix}) \\ & (E^{-(I * x)} - (2 * I) / E^{(3 * I * x)}) * \\ & (E^{(I * x)} + 2 * I * E^{(3 * I * x)}) \end{aligned}$$

$$(e^{-ix} - 2ie^{-3ix})(e^{ix} + 2ie^{3ix})$$

$$\text{Simplify}[(e^{-ix} - 2ie^{-3ix})(e^{ix} + 2ie^{3ix})]$$

$$5 - 2ie^{-2ix} + 2ie^{2ix}$$

$$\text{Simplify}[(e^{-ix} + 2ie^{-3ix})(e^{ix} + 2ie^{3ix})]$$

$$-3 + 2ie^{-2ix} + 2ie^{2ix}$$

ReleaseHold[HoldComplete[Null]]

$$\int_{-\pi}^{\pi} (5 - 2ie^{-2ix} + 2ie^{2ix}) dx$$

$$10\pi$$

(b) ψ can be seen as the superposition of the two waves: e^{ix} and e^{-3ix} . From the each wave, $\hat{p}_x e^{ix} = -i\hbar \frac{\partial}{\partial x} e^{ix} = \hbar e^{ix}$, $\hat{p}_x e^{-3ix} = -i\hbar \frac{\partial}{\partial x} e^{-3ix} = 3\hbar e^{-3ix}$. Therefore, the values of the momentum are \hbar and $3\hbar$. Since the ratio of coefficients of the two wavefunctions is 1:2, the ratio of the probability is 1:4. Therefore, the probability for the momentum \hbar is $\frac{1}{5}$, and the probability for $3\hbar$ is $\frac{4}{5}$.

(c) The probability of finding the particle is expressed as $\psi^* \psi = |\psi|^2$.

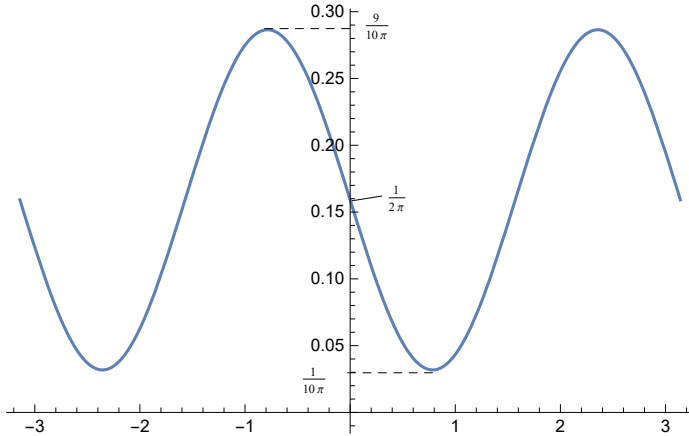
$$\psi^* \psi = \frac{1}{2\pi} + \frac{i}{5\pi}(e^{2ix} - e^{-2ix}) = \frac{1}{2\pi} - \frac{2}{5\pi} \sin 2x$$

$$\text{Simplify}\left[\frac{1}{\sqrt{10\pi}}(e^{-ix} - 2ie^{-3ix}) \frac{1}{\sqrt{10\pi}}(e^{ix} + 2ie^{3ix})\right]$$

$$\frac{(e^{-ix} - 2ie^{-3ix})(e^{ix} + 2ie^{3ix})}{10\pi}$$

(d) The graph is below:

$$\text{Plot}[1 / (2 * \text{Pi}) - (2 * \text{Sin}[2 * x]) / (5 * \text{Pi}), \{x, -\text{Pi}, \text{Pi}\}]$$



(e) The expectation value of the position is expressed as $\int_{-\pi}^{\pi} \psi^* x \psi dx$

$$\int_{-\pi}^{\pi} \psi^* x \psi dx = \int_{-\pi}^{\pi} x \left(\frac{1}{2\pi} - \frac{2}{5\pi} \sin 2x \right) dx = \frac{2}{5}$$

Integrate $\left[x \left(\frac{1}{2\pi} - \frac{2}{5\pi} \sin[2x] \right), \{x, -\pi, \pi\} \right]$

$$\frac{2}{5}$$

Problem 2. (20 points)

Show that the velocity of an electron in the n^{th} Bohr's orbit of the hydrogen atom is described as $V = \frac{e^2}{2 \epsilon_0 n \hbar}$. And calculate the velocity and de Broglie's wavelengths of electrons in the first three Bohr's orbits.

The photon wavelength λ is described as $\lambda = \frac{h}{p}$ with Planck's constant h and the momentum p . With this, de Broglie's wavelength is given by $\lambda = \frac{h}{mv}$ (1). Bohr's condition for orbit stability is $2\pi r = n\lambda$ (2), where r is the radius of the orbit that contain n wavelengths, and n is called the quantum number of the orbit. With eqs. (1) and (2), the angular momentum is quantized as $mvr = \frac{n\hbar}{2\pi} = n\hbar$ (3). From the balance between the coulomb force and the centrifugal force, we have $\frac{e^2}{4\pi\epsilon_0 r^2} = \frac{mv^2}{r}$ (4). With eqs. (3) and (4), r can be deduced to be $r = \frac{\epsilon_0 \hbar^2 n^2}{\pi m e^2}$ (5). Therefore, plugging-in (5) for (3), we have $v = \frac{n\hbar}{2\pi m r} = \frac{n\hbar \pi m e^2}{2\pi m \epsilon_0 \hbar^2 n^2} = \frac{e^2}{2 \epsilon_0 n \hbar}$ (6).

With eqs. (1) and (6), de Broglie's wavelength is given by $\lambda = \frac{h}{mv} = \frac{h 2 \epsilon_0 n \hbar}{m e^2} = \frac{2 \epsilon_0 n \hbar^2}{m e^2}$ (7).

The velocity and de Broglie's wavelength for $n=1$ are calculated as

$$v_1 = \frac{e^2}{2 \epsilon_0 \hbar} = \frac{(1.602 \times 10^{-19})^2}{2 \times 8.854 \times 10^{-12} \times 6.626 \times 10^{-34}} = 2.187 \times 10^6 \text{ m/s}$$

$$v_n = \frac{v_1}{n}, \text{ therefore, } v_2 = 1.094 \times 10^6 \text{ m/s, and } v_3 = 0.729 \times 10^6 \text{ m/s}$$

$$\lambda_1 = \frac{h}{mv_1} = \frac{6.626 \times 10^{-34}}{9.109 \times 10^{-31} \times 2.187 \times 10^6} = 0.3326 \times 10^{-9} \text{ m} = 333 \text{ pm}$$

$\lambda_n = n \lambda_1$, therefore, $\lambda_2 = 665 \text{ pm}$, and $\lambda_3 = 997 \text{ pm}$.

Problem 3. (20 points)

Consider one dimension and answer the questions below.

(a) If the wavefunction of a particle: $\psi(x) = \sqrt{a} e^{-a|x|}$ (a : the real and positive constant) is given, calculate the probability of finding the particle in the region of $|x| \leq L$.

(b) Defining two operators \hat{a} and \hat{b} as $\hat{a} = \frac{1}{\sqrt{2\hbar}}(\hat{p} - i\hat{x})$, $\hat{b} = \frac{1}{\sqrt{2\hbar}}(\hat{p} + i\hat{x})$, respectively, where \hat{x} and \hat{p} are the position and momentum operators, calculate the commutation relation $[\hat{a}, \hat{b}]$.

(a) The probability can be written as $\int_{-L}^L \psi^* \psi dx$

$$\int_{-L}^L \psi^* \psi dx = \int_{-L}^L (\sqrt{a})^2 e^{-2a|x|} dx = a \int_{-L}^0 e^{2ax} dx + a \int_0^L e^{-2ax} dx = 1 - e^{-2aL}$$

Integrate $[E^{\wedge}(2 * a * x), \{x, -L, 0\}] + \text{Integrate}[E^{\wedge}(-2 * a * x), \{x, 0, L\}]$

$$= \frac{-1 + e^{-2aL}}{a}$$

(b) $[\hat{a}, \hat{b}] = [\frac{1}{\sqrt{2\hbar}}(\hat{p} - i\hat{x}), \frac{1}{\sqrt{2\hbar}}(\hat{p} + i\hat{x})] = \frac{1}{2\hbar} \{[\hat{p}, \hat{p}] - [i\hat{x}, \hat{p}] + [\hat{p}, i\hat{x}] - [-i\hat{x}, i\hat{x}]\} = -\frac{i}{\hbar} [\hat{x}, \hat{p}] = -\frac{i}{\hbar} i\hbar = 1$

Problem 4. (30 points)

Consider a wavefunction $\psi(\vec{r}, t) = \sum_n c_n(t) \phi_n(\vec{r})$ with the orthonormalized functions $\phi_n(\vec{r})$ that satisfy $\hat{H}\phi_n = \hbar\omega_n\phi_n$ (n : natural number).

(a) Using the time-dependent Schrödinger equation, derive a derivative (differential) equation that the time-dependent coefficients $c_n(t)$ satisfy.

(Obtain a derivative (differential) equation of $c_n(t)$ from the time-dependent Schrödinger equation.)

(b) Obtain the wavefunction at time t . The initial conditions are $c_1(0) = \alpha$, $c_2(0) = \beta$ and others $c_n(0) = 0$ (α, β : the positive real numbers, and $\alpha^2 + \beta^2 = 1$)

(c) With the wavefunction obtained in (b), calculate the expectation value of the energy of that state.

(a) Plugging in $\psi(\vec{r}, t) = \sum_n c_n(t) \phi_n(\vec{r})$ for the time-dependent Schrödinger equation: $\hat{H}\psi = i\hbar \frac{\partial \psi}{\partial t}$,

$$\hat{H} \sum_n c_n(t) \phi_n(\vec{r}) = i\hbar \frac{\partial}{\partial t} \sum_n c_n(t) \phi_n(\vec{r})$$

$$\sum_n c_n(t) \hbar\omega_n \phi_n(\vec{r}) = i\hbar \sum_n \frac{d}{dt} c_n(t) \phi_n(\vec{r})$$

Multiplying the both sides by $\phi_k^*(\vec{r})$,

$$\sum_n c_n(t) \hbar\omega_n \int \phi_k^*(\vec{r}) \phi_n(\vec{r}) dV = i\hbar \sum_n \frac{d}{dt} c_n(t) \int \phi_k^*(\vec{r}) \phi_n(\vec{r}) dV.$$

Due to the orthogonality of the eigenfunctions $\int \phi_k^*(\vec{r}) \phi_n(\vec{r}) dV = \delta_{kn}$,

$$\sum_n c_n(t) \omega_n \delta_{kn} = i \sum_n \frac{d}{dt} c_n(t) \delta_{kn}$$

Therefore, the required derivative equation is expressed as

$$i \frac{d}{dt} c_n(t) = \omega_n c_n(t)$$

(b) The solution of the derivative equation obtained in (a) is given by

$$c_n(t) = c_n(0) e^{-i\omega_n t}$$

Since the initial conditions are $c_1(0) = \alpha$, $c_2(0) = \beta$ and others $c_n(0) = 0$, the wavefunction at time t can be expressed as

$$\psi(\vec{r}, t) = \alpha e^{-i\omega_1 t} \phi_1(\vec{r}) + \beta e^{-i\omega_2 t} \phi_2(\vec{r})$$

(c) The expectation value of the energy of that state is $\langle \hat{\mathcal{H}} \rangle$, and with the relation of $\int \phi_k^*(\vec{r}) \phi_n(\vec{r}) dV = \delta_{kn}$, it can be calculated as follows:

$$\begin{aligned} \langle \hat{\mathcal{H}} \rangle &= \int \psi^*(\vec{r}, t) \hat{\mathcal{H}} \psi(\vec{r}, t) dV = \\ &= \int \{ \alpha e^{-i\omega_1 t} \phi_1(\vec{r}) + \beta e^{-i\omega_2 t} \phi_2(\vec{r}) \}^* \hat{\mathcal{H}} \{ \alpha e^{-i\omega_1 t} \phi_1(\vec{r}) + \beta e^{-i\omega_2 t} \phi_2(\vec{r}) \} dV \\ &= \int \{ \alpha e^{i\omega_1 t} \phi_1^*(\vec{r}) + \beta e^{i\omega_2 t} \phi_2^*(\vec{r}) \} \{ \alpha e^{-i\omega_1 t} \hbar \omega_1 \phi_1(\vec{r}) + \beta e^{-i\omega_2 t} \hbar \omega_2 \phi_2(\vec{r}) \} dV \\ &= \alpha^2 \hbar \omega_1 + \beta^2 \hbar \omega_2 \end{aligned}$$

Problem 5. (20 points)

Assume the Hamiltonian of a particle in three dimension is expressed by spherical-polar coordinates ($x = r \sin\theta \cos\phi$, $y = r \sin\theta \sin\phi$, $z = r \cos\theta$) as below:

$$\hat{\mathcal{H}} = -\frac{\hbar^2}{2m} \left[\frac{d^2}{dr^2} + \frac{2}{r} \frac{d}{dr} + \frac{1}{r^2} \frac{\partial^2}{\partial \theta^2} + \frac{\cos\theta}{r^2 \sin\theta} \frac{\partial}{\partial \theta} + \frac{1}{r^2 \sin^2\theta} \frac{\partial^2}{\partial \phi^2} \right] - \frac{k}{r} \quad (k: \text{the positive real constant})$$

constant)

(a) If a wavefunction: $\psi = N e^{-br}$ (N, b : the positive real constant) is given, determine the value of b so that ψ is the eigenfunction of $\hat{\mathcal{H}}$ and also derive the eigenvalue.

(b) Normalize ψ and deduce N that is expressed by b .

(a) The wavefunction $\psi = N e^{-br}$ depends only on r .

$$\hat{\mathcal{H}}\psi = \left[-\frac{\hbar^2}{2m} \nabla^2 - \frac{k}{r} \right] \psi = -\frac{\hbar^2}{2m} \left[b^2 - \frac{2b}{r} \right] N e^{-br} - \frac{k}{r} N e^{-br} = -\frac{\hbar^2 b^2}{2m} \psi + \left[\frac{\hbar^2 b}{m} - k \right] \frac{1}{r} \psi.$$

Here, we use $\frac{d}{dr} e^{-br} = -b e^{-br}$, $\frac{d^2}{dr^2} e^{-br} = \frac{d}{dr} (-b e^{-br}) = b^2 e^{-br}$.

In order for ψ to be an eigenfunction of $\hat{\mathcal{H}}$ for any r ,

$$\frac{\hbar^2 b}{m} - k = 0$$

$$\therefore b = \frac{mk}{\hbar^2}$$

The eigenvalue for this,

$$-\frac{\hbar^2 b^2}{2m} = -\frac{\hbar^2}{2m} \frac{m^2 k^2}{\hbar^4} = -\frac{mk^2}{2\hbar^2}$$

(b) Plugging in $\psi = N e^{-br}$ for $\int \psi^* \psi dV = 1$, and using partial integral for $\int_0^\infty e^{-2br} r^2 dr$,

$$\int_0^\infty \int_0^\pi \int_0^{2\pi} N^2 e^{-2br} r^2 \sin\theta dr d\theta d\phi = 4\pi N^2 \int_0^\infty e^{-2br} r^2 dr = 4\pi N^2 \frac{1}{4b^3} = \frac{\pi N^2}{b^3} = 1$$

$$\text{Therefore, } \mathbf{N} = \sqrt{\frac{b^3}{\pi}}$$

Integrate $[r^2 \text{Exp}[-2 b r], \{r, 0, \infty\}]$

ConditionalExpression $\left[\frac{1}{4 b^3}, \text{Re}[b] > 0\right]$

Problem 6. (25 points)

An electron trapped in an infinite depth well of width $L=1\text{nm}$. Consider the transition from the excited state $n = 2$ to the ground state $n = 1$. Calculate the wavelength of light emitted.

Describe and explain all the processes of how to derive energy and wavelength emitted. (Key points and processes): Create Schrödinger's equation. Derive a general solution. Make clear boundary conditions. Discuss energy states.

The one-dimensional Schrödinger equation along the x direction is expressed as

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + V\psi = E\psi,$$

where ψ is the wavefunction of the particle, E is the total energy of the particle, and V is the potential energy. We consider the potential energy V is ∞ in the region of $x \leq 0$, and $x \geq L$, and V is 0 in the range of $0 < x < L$.

Therefore, in the region of $x \leq 0$, and $x \geq L$, $\psi(x) = 0$, and in the region of $0 < x < L$, The Schrödinger equation is rewritten by

$$-\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} = E\psi.$$

With $k = \frac{\sqrt{2mE}}{\hbar}$, $\frac{d^2\psi}{dx^2} = -k^2\psi$. The general solution of this differential equation is $\psi(x) = A\sin(kx) + B\cos(kx)$, where A and B are the constants. To determine A , B and k , we use the boundary conditions; those are $\psi(0) = 0$ and $\psi(L) = 0$. From the former, we obtain $B = 0$, and from the latter, $\psi(L) = A\sin(kL) = 0$. If $A=0$, $\psi(x) = 0$. This means no particle exists, which is unreasonable. Therefore, $A \neq 0$, and the possible values of kL are $kL = 0, \pm\pi, \pm2\pi, \dots$. However, $kL=0$ means $\psi(x)=0$ since E becomes 0, and the constant A absorbs the sign since $\sin(-x) = -\sin x$. Therefore, we obtain the condition as below: $kL = \pi, 2\pi, 3\pi, \dots = n\pi$ ($n = 1, 2, 3, \dots$)

We now obtain,

$$\frac{\sqrt{2mE}}{\hbar} = n\pi \quad (n = 1, 2, 3, \dots).$$

As a result,

$$E = \frac{\hbar^2 \pi^2 n^2}{2mL^2} \quad (n = 1, 2, 3, \dots)$$

This means the energy states are not continuous but discrete depending on natural number n .

Using the expression of the energy E , the transition from the excited state $n = 2$ to the

ground state $n = 1$ is given by

$$E_2 - E_1 = \frac{3\pi^2 \hbar^2}{2mL^2} = \frac{3\hbar^2}{8mL^2}$$

This energy gap corresponds to the emitted photon energy $h\nu$. Therefore,

$$h\nu = h \frac{c}{\lambda} = \frac{3\hbar^2}{8mL^2} = \frac{3 \times (6.626 \times 10^{-34})^2}{8 \times (9.1095 \times 10^{-31}) \times (10^{-9})^2} = 1.80734 \times 10^{-19} \text{ J}$$

The wavelength emitted is obtained as below:

$$\lambda = \frac{hc}{1.80734 \times 10^{-19}} = \frac{6.626 \times 10^{-34} \times 2.998 \times 10^8}{1.80734 \times 10^{-19}} = 1.09912 \times 10^{-6} \text{ m} = 1.1 \mu\text{m}$$

$$\frac{3 (6.626 \times 10^{-34})^2}{8 (9.1095 \times 10^{-31}) (10^{-9})^2}$$

$$1.80734 \times 10^{-19}$$

$$\frac{6.626 \times 10^{-34} \times 2.998 \times 10^8}{1.80734 \times 10^{-19}}$$

$$1.09912 \times 10^{-6}$$

* Physical constants and some formula:

Speed of light: $c = 2.998 \times 10^8 \text{ m/s} = 3.0 \times 10^8 \text{ m/s}$

Planck's constant: $h = 6.626 \times 10^{-34} \text{ J}\cdot\text{s} = 6.6 \times 10^{-34} \text{ J}\cdot\text{s}$

Electron charge: $e = 1.602 \times 10^{-19} \text{ C} = 1.6 \times 10^{-19} \text{ C}$

Electron rest mass: $m_e = 9.1095 \times 10^{-31} \text{ Kg} = 9.1 \times 10^{-31} \text{ Kg}$

Dielectric constant in vacuum: $\epsilon_0 = 8.854 \times 10^{-12} \text{ C}^2 \text{ N}^{-1} \text{ m}^{-2} = 8.9 \times 10^{-12} \text{ C}^2 \text{ N}^{-1} \text{ m}^{-2}$