Chapter 6 The Laplace Transform

6-1 Definition of the Laplace Transform

1. For the linear time-invariant system with impulse response h(t), the output y(t) corresponding to the input of the form e^{st} is

$$y(t) = \underbrace{H(s)}_{\text{eigenvalue}} \cdot \underbrace{e^{st}}_{\text{eigenfunction}}$$
(6.1)

where

$$H(s) = \int_{-\infty}^{\infty} h(t) e^{-st} dt$$
(6.2)

is referred to as the Laplace transform of h(t).

$$s = j\omega \Rightarrow H(j\omega) = \int_{-\infty}^{\infty} h(t)e^{-j\omega t}dt$$
 is the Fourier transform of $h(t)$.

2. The Laplace transform of a general signal x(t):

$$\begin{cases} X(s) \triangleq \int_{-\infty}^{\infty} x(t) e^{-st} dt = \mathcal{L} \{ x(t) \} \\ x(t) \longleftrightarrow X(s) \end{cases}$$
(6.3)

$$X(s)\Big|_{s=j\omega} = X(j\omega) = \mathcal{F}\{x(t)\}$$
(6.4)

3. The Laplace transform of x(t) can be interpreted as the Fourier transform of x(t) after multiplication by a real exponential. $s = \sigma + j\omega$

$$X(s) = X(\sigma + j\omega) = \int_{-\infty}^{\infty} x(t) e^{-(\sigma + j\omega)t} dt$$

=
$$\int_{-\infty}^{\infty} [x(t)e^{-\sigma t}] e^{-j\omega t} dt = \mathbf{\mathcal{F}} \{x(t)e^{-\sigma t}\}$$
 (6.5)

Example 6.1: $x(t) = e^{-at}u(t), X(j\omega)$ converges for $a > 0 \left(\int_{-\infty}^{\infty} |x(t)|^2 dt < \infty\right)$.

$$X(j\omega) = \mathcal{F}\left\{x(t)\right\} = \int_0^\infty e^{-at} e^{-j\omega t} dt = \frac{1}{j\omega + a}, \ a > 0$$
$$X(s) = \mathcal{L}\left\{x(t)\right\} = \int_0^\infty e^{-(s+a)t} dt$$

With $s = \sigma + j\omega$, we have

$$X(\sigma + j\omega) = \int_0^\infty e^{-(\sigma + a)t} \cdot e^{-j\omega t} dt = \mathcal{F}\left\{e^{-(\sigma + a)t}u(t)\right\}$$
$$= \frac{1}{j\omega + (\sigma + a)}, \ \sigma + a > 0, \text{ i.e., } \sigma > -a \text{ or } \operatorname{Re}\left\{s\right\} > -a$$
$$= \frac{1}{s + a}, \ \operatorname{Re}\left\{s\right\} > -a$$
$$\Rightarrow X(s) = \frac{1}{s + a}, \ \operatorname{Re}\left\{s\right\} > -a$$

Note:

- The Laplace transform converges for some values of Re{s}, and not for the others.
- The existence of the Laplace transform does not imply the existence of the Fourier transform, e.g., $x(t) = e^{-at}u(t)$, a < 0.

Example 6.2: $x(t) = -e^{-at}u(-t)$

$$X(s) = -\int_{-\infty}^{\infty} e^{-at} e^{-st} u(-t) dt = -\int_{-\infty}^{0} e^{-(s+a)t} dt = \frac{1}{s+a}$$
$$\begin{pmatrix} = -\int_{-\infty}^{\infty} e^{-(\sigma+a)t} e^{-j\omega t} u(-t) dt = -\mathbf{\mathcal{F}} \left\{ e^{-(\sigma+a)t} u(-t) \right\} \\ (\because t < 0, \ \because \sigma + a < 0 \Rightarrow \sigma < -a) \end{pmatrix}$$

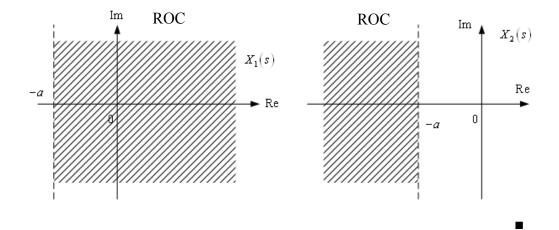
 $\Rightarrow X(s)$ exists only for $\operatorname{Re}\{s\} < -a$.

Note:

- In specifying the Laplace transform of a signal, both the algebraic expression and the range of values for which this expression is valid are required.
- The range of values for which the Laplace transform exists is referred to as the region of convergence (ROC) of the Laplace transform.

Example 6.3: Region of convergence (ROC) of $X_1(s) = \mathcal{L}\left\{e^{-at}u(t)\right\} = \frac{1}{s+a}$

and $X_2(s) = \mathcal{L}\{-e^{-at}u(-t)\} = \frac{1}{s+a}.$



Example 6.4: $x(t) = e^{-t}u(t) + e^{-2t}u(t)$

$$X(s) = \int_{-\infty}^{\infty} \left[e^{-t}u(t) + e^{-2t}u(t) \right] e^{-st} dt$$

$$= \int_{-\infty}^{\infty} e^{-t}e^{-st}u(t) dt + \int_{-\infty}^{\infty} e^{-2t}e^{-st}u(t) dt$$

$$= \frac{1}{s+1} + \frac{1}{s+2}, \quad \operatorname{Re}\{s\} > -1$$

$$\left\{ \underbrace{= \mathcal{L}\left\{e^{-t}u(t)\right\}}_{\frac{1}{s+1}, \operatorname{Re}\{s\} > -1} + \underbrace{\mathcal{L}\left\{e^{-2t}u(t)\right\}}_{\frac{1}{s+2}, \operatorname{Re}\{s\} > -2} \right\}$$

$$\Rightarrow \mathcal{L}\left\{e^{-t}u(t) + e^{-2t}u(t)\right\} = \frac{1}{s+1} + \frac{1}{s+2}, \quad \operatorname{Re}\{s\} > -1$$

$$= \frac{2s+3}{s^2+3s+2}, \quad \operatorname{Re}\{s\} > -1$$

$$= \frac{N(s)}{D(s)} \rightarrow \operatorname{numerator polynomial}$$

Note:

- Whenever x(t) is a linear combination of real or complex exponentials, X(s) can be expressed by X(s) = N(s)/D(s), i.e., X(s) is rational.
- The roots of the numerator polynomial (denominator polynomial) are referred to as the zeros (poles) of X(s) since for those values of s,

$$X(s) = 0 (X(s) \to \infty).$$

• X(s) = N(s)/D(s)

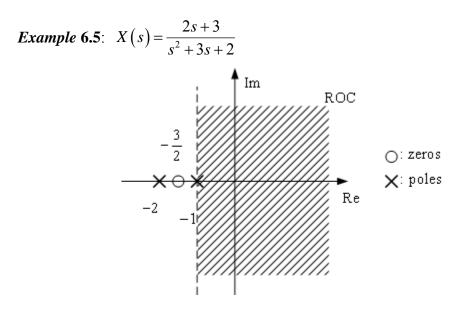
[The order of N(s)]<[The order of D(s)]

 \Rightarrow exist zeros at infinity $(s \rightarrow \infty, X(s) \rightarrow 0)$

[The order of N(s)]>[The order of D(s)]

 \Rightarrow exist poles at infinity $(s \rightarrow \infty, X(s) \rightarrow \infty)$

• The representation of X(s) = N(s)/D(s) through its poles and zeros in the *s*-plane is referred to as the pole-zero diagram or the pole-zero plot.



Example 6.6:
$$x(t) = \delta(t) - \frac{4}{3}e^{-t}u(t) + \frac{1}{3}e^{2t}u(t)$$

$$\delta(t) \xleftarrow{\mathcal{L}} 1, \text{ ROC : entire s plane}$$

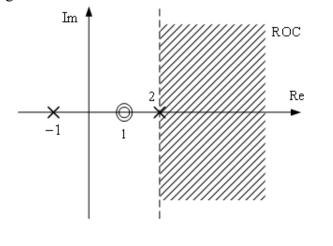
$$\frac{4}{3} e^{-t} u(t) \xleftarrow{\mathcal{L}} \frac{4}{3} \cdot \frac{1}{s+1}, \text{ Re}\{s\} > -1$$

$$\frac{1}{3} e^{2t} u(t) \xleftarrow{\mathcal{L}} \frac{1}{3} \cdot \frac{1}{s-2}, \text{ Re}\{s\} > 2$$

$$\left(\int_{-\infty}^{\infty} e^{2t} e^{-st} u(t) dt = \int_{-\infty}^{\infty} e^{(2-\sigma)t} e^{-j\omega t} u(t) dt, 2-\sigma < 0 \Rightarrow \sigma > 2 \Rightarrow \text{Re}\{s\} > 2\right)$$

$$X(s) = 1 - \frac{4}{3} \cdot \frac{1}{s+1} + \frac{1}{3} \cdot \frac{1}{s-2}, \quad \operatorname{Re}\{s\} > 2$$
$$= \frac{(s-1)^2}{(s+1)(s-2)}, \quad \operatorname{Re}\{s\} > 2$$

We will refer to the order of pole or zero as the number of times it is repeated at a given location.



6-2 The Region of Convergence for Laplace Transforms

Properties of ROC for Laplace transforms:

- The ROC of X(s) consists of strips parallel to the jω-axis in the s-plane.
 Example 6.7: X(s) converges only for Re{s} > a (or Re{s} < a). The ROC depends only on the real part of s.
- 2. For rational Laplace transforms, the ROC does not contain any poles.

$$s = \text{pole} \implies X(s) \rightarrow \infty$$

3. If x(t) is of finite duration and if there is at least one value of *s* for which the Laplace transform converges, then the ROC is the entire *s*-plane.



Figure 6.1 Finite-duration signal.

Proof:

Let x(t) be zero outside the interval between T_1 and T_2 . Then

$$X(s) = \int_{T_1}^{T_2} x(t) e^{-st} dt$$
 (6.6)

Assume the line $\operatorname{Re}\{s\} = \sigma_0$ is in the ROC. Then

$$\int_{T_1}^{T_2} \left| x(t) \right| e^{-\sigma_0 t} dt < \infty \tag{6.7}$$

(1) For $\sigma_1 > \sigma_0$,

$$\int_{T_{1}}^{T_{2}} |x(t)| e^{-\sigma_{1}t} dt = \int_{T_{1}}^{T_{2}} |x(t)| e^{-\sigma_{0}t} e^{-(\sigma_{1}-\sigma_{0})t} dt$$

$$< e^{-(\sigma_{1}-\sigma_{0})T_{1}} \int_{T_{1}}^{T_{2}} |x(t)| e^{-\sigma_{0}t} dt < \infty$$
(6.8)

(: The maximum value of $e^{-(\sigma_1 - \sigma_0)t}$ in the interval of integration is $e^{-(\sigma_1 - \sigma_0)T_1}$.)

This implies the *s*-plane for $\operatorname{Re}\{s\} > \sigma_0$ is in the ROC.

(2) For
$$\sigma_2 < \sigma_0$$
,

$$\int_{T_{1}}^{T_{2}} |x(t)| e^{-\sigma_{2}t} dt = \int_{T_{1}}^{T_{2}} |x(t)| e^{-\sigma_{0}t} e^{-(\sigma_{2}-\sigma_{0})t} dt$$

$$< e^{-(\sigma_{2}-\sigma_{0})T_{2}} \int_{T_{1}}^{T_{2}} |x(t)| e^{-\sigma_{0}t} dt < \infty$$
(6.9)

This implies the *s*-plane for $\operatorname{Re}\{s\} < \sigma_0$ is in the ROC.

From Eq. (6.1) and Eq. (6.2), we can see that the ROC of a finite-duration signal includes the entire *s*-plane.

4. If x(t) is right-sided and if the line $\operatorname{Re}\{s\} = \sigma_0$ is in the ROC, then all values of *s* for which $\operatorname{Re}\{s\} > \sigma_0$ will also be in the ROC.

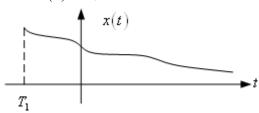


Figure 6.2 Right-sided signal.

 $x(t) = e^{t^2}u(t)$: there is no value of *s* for which the Laplace transform will converge.

Suppose the Laplace transform of x(t) converges for some value of σ , denoted by σ_0 . Then

$$\int_{-\infty}^{\infty} |x(t)| e^{-\sigma_0 t} dt < \infty$$
(6.10)

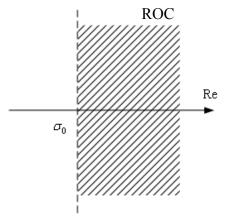
$$\Rightarrow \int_{T_1}^{\infty} |x(t)| e^{-\sigma_0 t} dt < \infty$$
(6.11)

For
$$\sigma_1 > \sigma_0$$
,

$$\Rightarrow \int_{T_1}^{\infty} |x(t)| e^{-\sigma_1 t} dt = \int_{T_1}^{\infty} |x(t)| e^{-\sigma_0 t} e^{-(\sigma_1 - \sigma_0)t} dt < e^{-(\sigma_1 - \sigma_0)T_1} \int_{T_1}^{\infty} |x(t)| e^{-\sigma_0 t} dt < \infty \quad (6.12)$$

(: The maximum value of $e^{-(\sigma_1 - \sigma_0)t}$ in the interval of integration is $e^{-(\sigma_1 - \sigma_0)T_1}$.)

ROC of a right-sided signal:



■ Figure 6.3 ROC of a right-sided signal.

5. If x(t) is left-sided and if the line $\operatorname{Re}\{s\} = \sigma_0$ is in the ROC, then all values of *s* for which $\operatorname{Re}\{s\} < \sigma_0$ will also be in the ROC. ROC of a left-sided signal:

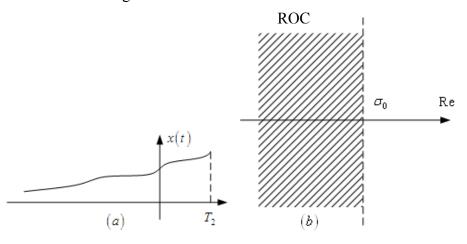


Figure 6.4 (a) Left-sided signal; (b) ROC of a left-sided signal.

6. If x(t) is two-sided and if the line $\operatorname{Re}\{s\} = \sigma_0$ is in the ROC, then the ROC will consist of a strip in the *s*-plane which includes the line $\operatorname{Re}\{s\} = \sigma_0$.

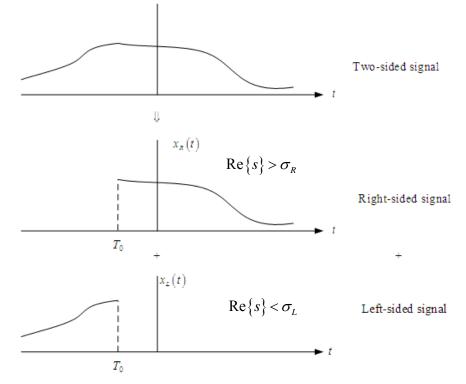


Figure 6.5 Two-sided signal divided into the sum of a right-sided and left-sided signal.

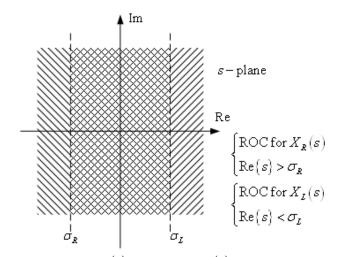


Figure 6.6 ROCs for $x_R(t)$ and for $x_L(t)$ assuming that they overlap. The overlap of the two ROCs is the ROC for $x(t) = x_R(t) + x_L(t)$. Note:

• σ_L must be greater than σ_R ; otherwise, the Laplace transform of x(t) does not exist.

Example 6.8:

$$x(t) = \begin{cases} e^{-at} , 0 < t < T \\ 0 , \text{ otherwise} \end{cases}$$

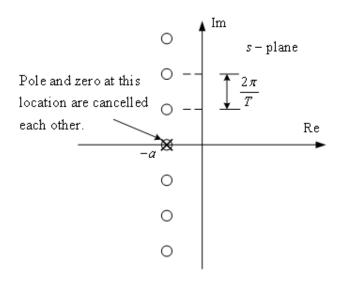
x(t) is a finite-duration sequence.

 \Rightarrow the ROC of X(s) in the entire *s*-plane

$$X(s) = \int_{0}^{T} e^{-at} e^{-st} dt = \frac{1}{s+a} \left[1 - e^{-(s+a)T} \right]$$
$$\begin{pmatrix} s = -a \Rightarrow 0 \\ 1 - e^{-(s+a)T} \to 0 \\ \vdots \\ \vdots \\ s \to -a \end{pmatrix}$$
$$\frac{d}{1 - e^{-(s+a)T}} = \lim_{s \to -a} \frac{d}{\frac{ds}{s}} \left[1 - e^{-(s+a)T} \right]$$
$$= \lim_{s \to -a} T e^{-aT} e^{-sT} = T$$

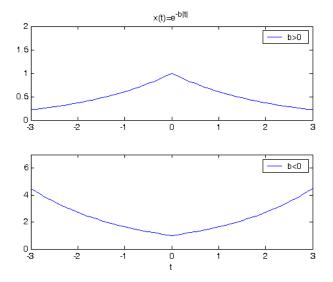
- $\begin{cases} X(s) \text{ has no poles.} \\ X(s) \text{ has an infinite number of zeros.} \end{cases}$

$$1 - e^{-(s+a)T} = 0 \Longrightarrow (s+a)T = j2\pi k, \ k = 0, \ \pm 1, \ \pm 2, \dots$$
$$\implies s = -a + j\frac{2\pi k}{T}, \ k = 0, \ \pm 1, \ \pm 2, \dots$$



Example 6.9: $x(t) = e^{-b|t|}$

$$x(t) = \underbrace{e^{-bt}u(t)}_{\text{right-sided}} + \underbrace{e^{bt}u(-t)}_{\text{left-sided}}$$
$$e^{-bt}u(t) \xleftarrow{\mathcal{L}} \frac{1}{s+b}, \quad \text{Re}\{s\} > -b$$
$$e^{bt}u(-t) \xleftarrow{\mathcal{L}} \frac{-1}{s-b}, \quad \text{Re}\{s\} < +b$$

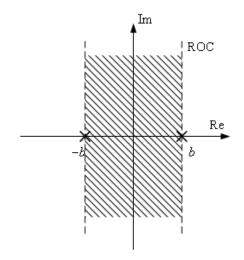


For b<0, there is no common region of convergence. $\Rightarrow x(t)$ has no Laplace transform if b<0.

For *b*> 0,

$$x(t) \longleftrightarrow \frac{1}{s+b} - \frac{1}{s-b} = \frac{-2b}{s^2 - b^2}, \quad -b < \operatorname{Re}\{s\} < b$$

Pole-zero plot



Fall 2011

Summary for the ROC:

Finite-duration signal
$$\rightarrow \begin{cases} \text{entire } s \text{-plane} \\ \text{does not exist} \end{cases}$$

Right-sided signal
$$\rightarrow \begin{cases} right-half s-plane \\ does not exist \end{cases}$$

Left-sided signal
$$\rightarrow \begin{cases} \text{left-half } s\text{-plane} \\ \text{does not exist} \end{cases}$$

Two-sided signal $\rightarrow \begin{cases} a \text{ strip} \\ does \text{ not exist} \end{cases}$

Note:

- The ROC is bounded by poles or extends to infinity.
- For a right-sided signal, the ROC is the region in the *s*-plane to the right of the rightmost pole.
- For a left-sided signal, the ROC is the region in the *s*-plane to the left of the leftmost pole.

6-3 The Inverse Laplace Transform

$$X(\sigma + j\omega) = \int_{-\infty}^{\infty} (x(t)e^{-\sigma t})e^{-j\omega t} dt = \mathbf{\mathcal{T}}\left\{x(t)e^{-\sigma t}\right\}$$
(6.13)

$$\Rightarrow x(t)e^{-\sigma t} = \mathbf{\mathcal{F}}^{-1}\left\{X\left(\sigma+j\omega\right)\right\} = \frac{1}{2\pi}\int_{-\infty}^{\infty}X\left(\sigma+j\omega\right)e^{j\omega t}d\omega \qquad (6.14)$$

$$\Rightarrow x(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(\sigma + j\omega) e^{j\omega t} e^{\sigma t} d\omega = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(\sigma + j\omega) e^{(\sigma + j\omega)t} d\omega \quad (6.15)$$

If we change the variable of integration from ω to s and use the fact that σ is constant so that $ds = jd\omega$, we obtain

$$x(t) = \frac{1}{2\pi j} \int_{\sigma - j\infty}^{\sigma + j\infty} X(s) e^{st} ds$$
(6.16)

"The basic Inverse Laplace Transform." Note:

• σ is any value in the ROC of X(s).

Example 6.10:

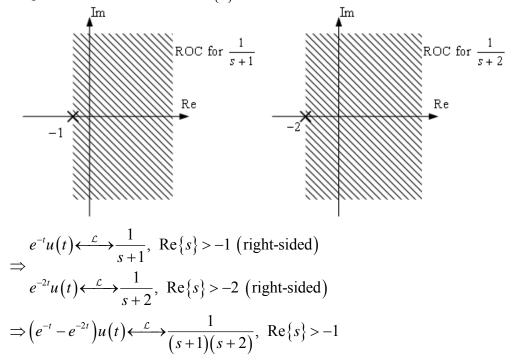
$$X(s) = \frac{1}{(s+1)(s+2)}, \text{ Re}\{s\} > -1$$

$$X(s) = \frac{1}{(s+1)(s+2)} = \frac{A}{s+1} + \frac{B}{s+2} = \frac{1}{s+1} - \frac{1}{s+2}$$

$$A = (s+1)X(s)|_{s=-1} = 1$$

$$B = (s+2)X(s)|_{s=-2} = -1$$

Since the ROC for X(s) is $\operatorname{Re}\{s\} > -1$, the ROC for the individual terms in the partial fraction includes $\operatorname{Re}\{s\} > -1$.



Example 6.11:

$$X(s) = \frac{1}{(s+1)(s+2)}, \operatorname{Re}\{s\} < -2 \text{ (left-sided)}$$
$$= \frac{1}{(s+1)(s+2)} = \frac{1}{s+1} - \frac{1}{s+2}$$
$$x(t) = (-e^{-t} + e^{-2t})u(-t) \longleftrightarrow \frac{2}{(s+1)(s+2)}, \operatorname{Re}\{s\} < -2$$
$$\because e^{-bt}u(-t) \longleftrightarrow \frac{2}{s-b}, \operatorname{Re}\{s\} < b$$
$$\vdots \begin{cases} e^{-t}u(-t) \longleftrightarrow \frac{2}{s+1}, \operatorname{Re}\{s\} < -1\\ \vdots \end{cases} \begin{cases} e^{-2t}u(-t) \longleftrightarrow \frac{2}{s+2}, \operatorname{Re}\{s\} < -2 \end{cases}$$

6-4 Geometric Evaluation of the Fourier Transform from the Pole-Zero Plot

1. Consider $H(s_1) = s_1 - a$

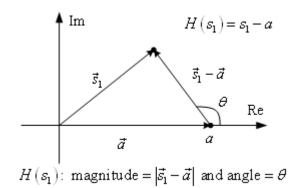


Figure 6.7 Complex plane representation of the vectors \vec{s}_1 , \vec{a} , and $(\vec{s}_1 - \vec{a})$ representing the complex numbers s_1 , a, and $(s_1 - a)$ respectively.

For H(s)=1/(s-a), the denominator can be represented by the same vector as above and the value of $H(s_1)$ has a magnitude that is the reciprocal of the vector $(\vec{s}_1 - \vec{a})$.

2. A system function described by a linear differential equation with constant coefficients is a rational fraction of the form

$$H(s) = \frac{B(s)}{A(s)} = \frac{\sum_{k=0}^{M} b_k s^k}{\sum_{k=0}^{N} a_k s^k}.$$
(6.17)

The numerator and denominator polynomials B(s) and A(s) can always be factored into products of M and N first-order terms, respectively, and thus H(s) can also be written as

$$H(s) = C \frac{\prod_{k=1}^{M} (s - z_k)}{\prod_{k=1}^{N} (s - p_k)},$$
(6.18)

where $C = b_M / a_N$, and z_k and p_k are the zeros and poles, respectively, of H(s). Therefore the corresponding frequency response $H(j\omega)$ is simply

$$H(j\omega) = C \frac{\prod_{k=1}^{M} (j\omega - z_k)}{\prod_{k=1}^{N} (j\omega - p_k)}.$$
(6.19)

For a given frequency ω , each complex-valued numerator term $(j\omega - z_k)$ in Eq. (6.19) can be thought of as a vector in the complex (s) plane from the zero z_k to the point $j\omega$ on the imaginary axis; and likewise, each denominator term $(j\omega - p_k)$ is effectively a vector from the pole p_k to the point $j\omega$. Hence, via Eq. (6.19), the magnitude and phase responses of the system can be determined by the lengths and angles, respectively, of these pole/zero vectors as functions of the variable ω .

To evaluate H(s) at $s = s_1$, each term in the product is represented by a vector from the zero or pole to the point s_1 :

(1) The magnitude response $|H(j\omega)|$ is

$$\frac{|C|(\text{the product of the lengths of the zero vectors})}{(\text{the product of the lengths of the pole vectors})}$$

$$= |C| \frac{\prod_{k=1}^{M} |j\omega - z_{k}|}{\prod_{k=1}^{N} |j\omega - p_{k}|}.$$
(6.20)

The phase response $\angle H(j\omega)$ is

(the sum of the angles of the zero vectors)

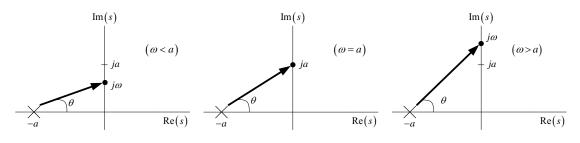
-(the sum of the angles of the pole vectors) (6.21)

$$= \sum_{k=1}^{M} \angle (j\omega - z_k) - \sum_{k=1}^{n} \angle (j\omega - p_k).$$

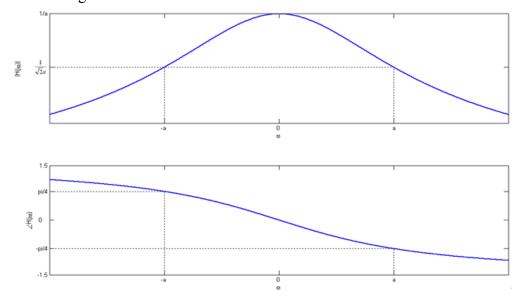
If C is negative, then in additional angle of π would be included.

In many cases, this simple geometric approach is sufficient to enable us to sketch $|H(j\omega)|$ and $\angle H(j\omega)$ with adequate directly from the pole/zero diagram for H(s) without having to evaluate $H(j\omega)$ itself.

The corresponding pole/zero plot showing the denominator vector $(j\omega + a)$ is drawn in the following with three cases:



The resulting sketch of $|H(j\omega)|$ and $\angle H(j\omega)$ are shown in the following.



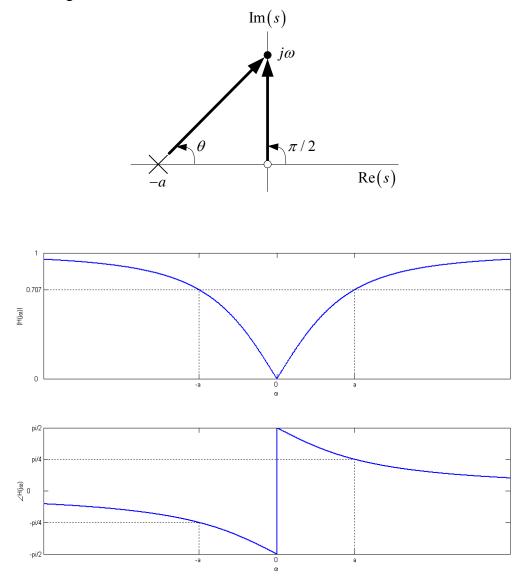
Note:

• $|j\omega + a|$ increases monotonically as increases from $\omega = 0$.

• The 3dB point will be at $|\omega| = a$.

• $\angle H(j\omega)$ decreases monotonically with ω from a value at $\omega = 0$ and approaches an asymptotic value of $-\pi/2$ for $\omega >> a$. Moreover, $\angle H(j\omega)$ is an odd function of ω .

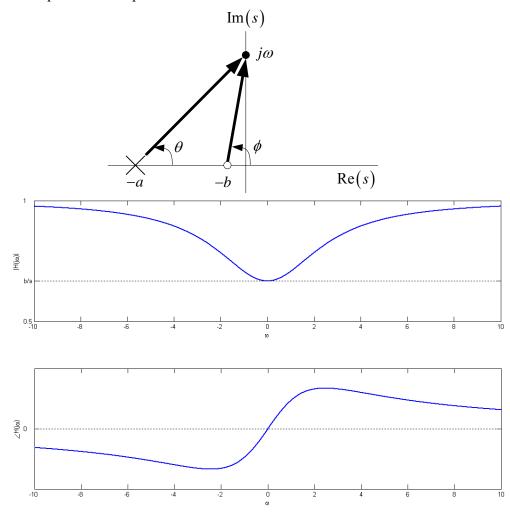
Example 6.14: Adding a zero at s = 0 to the preceding example, we have H(s) = s/(s+a), a > 0, and thus $H(j\omega) = j\omega/(j\omega+a)$. The corresponding pole/zero diagram, magnitude response, and phase response are shown in the following.



Note:

- $|H(j\omega)|$ is zero at $\omega = 0$ and decreases as $|\omega|$ increases. For $\omega >> a$, $|H(j\omega)| \approx 1$.
- The 3dB point will be at $\omega = a$.
- $\angle H(j\omega)$ discontinuity at $\omega = 0$ due to there is a zero on the $j\omega$ axis.

Example 6.15: Letting H(s) = (s+b)/(s+a), with 0 < b < a, we have the frequency response $H(j\omega) = (j\omega+b)/(j\omega+a)$. This case is almost the same as Example 6.14 except that the zero is moved to the left.



Note:

• $|H(j\omega)| = b/a < 1$ at $\omega = 0$ and increases as $|\omega|$ increases. For $\omega >> a$, $|H(j\omega)| \approx 1$. Therefore $|H(j\omega)|$ is again high-pass, but with less attenuation near $\omega = 0$.

6-5 Properties of the Laplace Transform

1. Linearity

$$x_1(t) \xleftarrow{\mathcal{L}} X_1(s)$$
, with ROC = R_1 (6.22)

$$x_2(t) \xleftarrow{\mathcal{L}} X_2(s)$$
, with ROC = R_2 (6.23)

 $\Rightarrow ax_1(t) + bx_2(t) \xleftarrow{\mathcal{L}} aX_1(s) + bX_2(s) \text{ with ROC containing } R_1 \cap R_2 \quad (6.24)$ Note:

• The ROC can also be larger than $R_1 \cap R_2$.

Example 6.16:
$$X_1(s) = \frac{1}{s+1}$$
, $\operatorname{Re}\{s\} > -1$
 $X_2(s) = \frac{1}{(s+1)(s+2)}$, $\operatorname{Re}\{s\} > -1$
 $x(t) = x_1(t) - x_2(t)$
 $X(s) = X_1(s) - X_2(s) = \frac{1}{s+1} - \frac{1}{(s+1)(s+2)}$
 $= \frac{s+1}{(s+1)(s+2)} = \frac{1}{s+2}$, $\operatorname{Re}\{s\} > -2$

In the combination of $x_1(t)$ and $x_2(t)$, the pole at s = -1 is cancelled by a zero at s = -1. \Rightarrow "pole-zero cancellation"

2. Time shifting

$$x(t) \xleftarrow{\mathcal{L}} X(s)$$
, with ROC = R (6.25)

$$\Rightarrow x(t-t_0) \xleftarrow{\mathcal{L}} e^{-st_0} X(s), \text{ with ROC} = R$$

$$\left(\int_{-\infty}^{\infty} x(t-t_0) e^{-st} dt = e^{-st_0} X(s) \right)$$
(6.26)

3. Shifting in the *s*-plane

$$x(t) \longleftrightarrow X(s)$$
, with ROC = R (6.27)

$$\Rightarrow e^{s_0 t} x(t) \longleftrightarrow X(s - s_0), \text{ with ROC} = R + \operatorname{Re}\{s_0\}$$

(:: pole $s_n \to s_n + s_0$) (6.28)

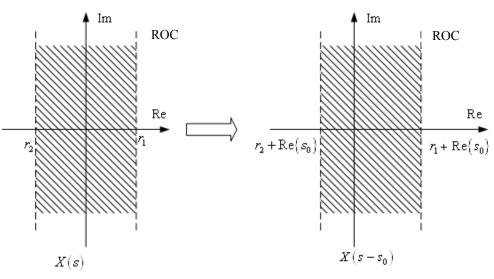


Figure 6.8 Effects on the ROC of shifting in the *s*-domain.

4. Time scaling

$$x(t) \xleftarrow{\mathcal{L}} X(s)$$
, with ROC = R (6.29)

$$\Rightarrow x(at) \longleftrightarrow \frac{1}{|a|} X\left(\frac{s}{a}\right), \text{ with ROC} = aR$$

$$(:: \text{ pole } s_p \to as_p)$$

$$(:: \text{ pole } s_p \to as_p)$$

$$(:: \text{ pole } s_p \to as_p)$$

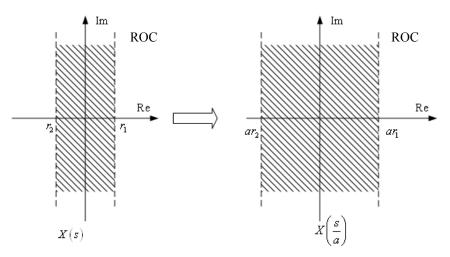


Figure 6.9 Effects on the ROC of time scaling.

a>0,

$$\mathcal{L}\left\{x(at)\right\} = \int_{-\infty}^{\infty} x(at)e^{-st}dt = \int_{-\infty}^{\infty} x(t')e^{-\frac{s-t'}{a}}\frac{1}{a}dt'$$
$$= \frac{1}{a}X\left(\frac{s}{a}\right) = \frac{1}{|a|}X\left(\frac{s}{a}\right)$$
(6.31)

a< 0,

$$\mathcal{L}\{x(at)\} = \int_{-\infty}^{\infty} x(at)e^{-st}dt = \int_{-\infty}^{\infty} x(t')e^{-\frac{s}{a}t'}\frac{1}{a}dt'$$

= $-\frac{1}{a}\int_{-\infty}^{\infty} x(t')e^{-\frac{s}{a}t'}dt' = \frac{1}{|a|}X\left(\frac{s}{a}\right)$ (6.32)

5. Convolution property

$$x_1(t) \xleftarrow{\mathcal{L}} X_1(s), \text{ with ROC} = R_1$$
 (6.33)

$$x_2(t) \xleftarrow{\mathcal{L}} X_2(s)$$
, with ROC = R_2 (6.34)

$$\Rightarrow x(t) = x_1(t) * x_2(t) \longleftrightarrow X(s) = X_1(s) X_2(s), \quad \text{with ROC} \\ \text{containing } R_1 \cap R_2 \quad (6.35)$$

Note:

• The ROC of X(s) may be larger than $R_1 \cap R_2$ if pole-zero cancellation occurs in the product.

Example 6.17:

$$X_1(s) = \frac{s+1}{s+2}, \text{ Re}\{s\} > -2$$
$$X_2(s) = \frac{s+2}{s+1}, \text{ Re}\{s\} > -1$$

Then $X(s) = X_1(s)X_2(s) = 1$, with ROC = entire *s*-plane.

6. Differentiation in the time domain

$$x(t) \longleftrightarrow X(s)$$
, with ROC = R (6.36)

$$\frac{dx(t)}{dt} \xleftarrow{\mathcal{L}} sX(s), \text{ with ROC containing } R$$
(6.37)

$$x(t) = \frac{1}{2\pi j} \int_{\sigma - j\infty}^{\sigma + j\infty} X(s) e^{st} ds \Longrightarrow \frac{d}{dt} x(t) = \frac{1}{2\pi j} \int_{\sigma - j\infty}^{\sigma + j\infty} sX(s) e^{st} ds \qquad (6.38)$$

Example 6.18:
$$x(t) = \frac{d^2}{dt^2} \left(e^{-3(t-2)} u(t-2) \right)$$

 $e^{-3t} u(t) \xleftarrow{\mathcal{L}} \frac{1}{s+3}$, with ROC Re $\{s\} > -3$
 $e^{-3(t-2)} u(t-2) \xleftarrow{\mathcal{L}} \frac{1}{s+3} e^{-2s}$, with ROC Re $\{s\} > -3$
 $x(t) = \frac{d^2}{dt^2} \left(e^{-3(t-2)} u(t-2) \right) \xleftarrow{\mathcal{L}} X(s) = \frac{s^2}{s+3} e^{-2s}$, with ROC Re $\{s\} > -3$

Example 6.19:
$$X(s) = \frac{2s^3 - 9s^2 + 4s + 10}{s^2 - 3s - 4}$$
, with $\operatorname{Re}\{s\} < -1$

$$\frac{2s - 3}{s^2 - 3s - 4}\overline{)2s^3 - 9s^2 + 4s + 10}$$

$$\frac{2s^3 - 6s^2 - 8s}{-3s^2 + 12s + 10}$$

$$\frac{-3s^2 + 9s + 12}{3s - 2}$$

$$X(s) = 2s - 3 + \frac{1}{s + 1} + \frac{2}{s - 4}$$
, with $\operatorname{Re}\{s\} < -1$

$$x(t) = 2\delta^{(1)}(t) - 3\delta(t) - e^{-t}u(-t) - 2e^{4t}u(-t)$$

Note:

- The ROC of sX(s) includes the ROC of X(s) and may be larger if X(s) has a first order pole at s = 0 which is cancelled by the multiplication by s.
- 7. Differentiation in the *s*-domain

$$x(t) \longleftrightarrow X(s)$$
, with ROC = R (6.39)

$$\Rightarrow -tx(t) \longleftrightarrow \frac{dX(s)}{ds}, \text{ with ROC} = R$$
(6.40)

$$X(s) = \int_{-\infty}^{\infty} x(t) e^{-st} dt$$
(6.41)

$$\frac{dX(s)}{ds} = \int_{-\infty}^{\infty} \left[-tx(t) \right] e^{-st} dt$$
(6.42)

8. Integration in the time domain

$$x(t) \xleftarrow{\mathcal{L}} X(s)$$
, with ROC = R (6.43)

$$\Rightarrow \int_{-\infty}^{t} x(\tau) d\tau \longleftrightarrow X(s)/s, \text{ with ROC containing } R \cap \{\operatorname{Re}\{s\} > 0\}$$
(6.44)

$$\int_{-\infty}^{t} x(\tau) d\tau = u(t) * x(t)$$
(6.45)

$$u(t) \xleftarrow{\mathcal{L}} s^{-1}, \text{ with ROC} = \operatorname{Re}\{s\} > 0$$

$$\begin{pmatrix} e^{-at}u(t) \xleftarrow{\mathcal{L}} \frac{1}{s+a}, \text{ with ROC} = \operatorname{Re}\{s\} > -a \\ \operatorname{Let} a = 0. \end{pmatrix}$$
(6.46)

$$x(t) \xleftarrow{\mathcal{L}} X(s)$$
, with ROC = R (6.47)

$$\Rightarrow \int_{-\infty}^{t} x(\tau) d\tau \longleftrightarrow \frac{1}{s} X(s), \text{ with ROC containing } R \cap \{ \operatorname{Re}\{s\} > 0 \}$$
(6.48)

9. The initial and final value theorems

x(t) = 0 for t < 0 and x(t) contains no impulses or higher-order singularities at the origin.

$$\begin{cases} x(0^{+}) = \lim_{s \to \infty} sX(s) \cdots \text{ The initial value theorem} \\ \lim_{t \to \infty} x(t) = \lim_{s \to 0} sX(s) \cdots \text{ The final value theorem} \end{cases}$$
(6.49)

Proof:

Expanding x(t) as a Taylor series at $t = 0^+$,

$$x(t) = \left[x(0^{+}) + x^{(1)}(0^{+})t + \dots + x^{(n)}(0^{+})\frac{t^{n}}{n!} + \dots \right] u(t)$$
 (6.50)

where $x^{(n)}(0^+)$ denotes the *n*th derivative of x(t) evaluated at $t = 0^+$.

$$u(t) \xleftarrow{\mathcal{L}} \frac{1}{s} \tag{6.51}$$

$$\begin{array}{c} tu(t) \xleftarrow{\mathcal{L}} \frac{1}{s^2} \\ \vdots \end{array} \tag{6.52}$$

$$\frac{t^n}{n!}u(t) \longleftrightarrow \frac{1}{s^n}$$
(6.53)

$$\Rightarrow \mathcal{L}\left\{x(t)\right\} = \frac{1}{s}x(0^{+}) + \frac{1}{s^{2}}x^{(1)}(0^{+}) + \dots + \frac{1}{s^{n}}x^{(n)}(0^{+}) + \dots = X(s)$$
(6.54)

$$\Rightarrow sX(s) = x(0^{+}) + \frac{1}{s}x^{(1)}(0^{+}) + \dots + \frac{1}{s^{n-1}}x^{(n)}(0^{+}) + \dots$$
(6.55)

$$\Rightarrow \lim_{s \to \infty} sX(s) = x(0^+) \cdots \cdots \text{The initial value theorem}$$
(6.56)

Let us consider the limit of the integral $\int_{0^+}^{\infty} \frac{dx(t)}{dt} e^{-st} dt$ as s approach 0. We have

$$\lim_{s \to 0} \int_{0^+}^{\infty} \frac{dx(t)}{dt} e^{-st} dt = \int_{0^+}^{\infty} \frac{dx(t)}{dt} dt = x(t) \Big|_{0^+}^{\infty} = \lim_{t \to \infty} x(t) - x(0^+)$$
(6.57)

Also,

$$\lim_{s \to 0} \int_{0^{+}}^{\infty} \frac{dx(t)}{dt} e^{-st} dt = \lim_{s \to 0} \left\{ \left[x(t) e^{-st} \right]_{0^{+}}^{\infty} - \int_{0^{+}}^{\infty} x(t) \frac{de^{-st}}{dt} dt \right\}$$
$$= \lim_{s \to 0} \left\{ \left[x(t) e^{-st} \right]_{0^{+}}^{\infty} - \int_{0^{+}}^{\infty} x(t) (-s) e^{-st} dt \right\}$$
$$= \lim_{s \to 0} \left[-x(0^{+}) + sX(s) \right] = -x(0^{+}) + \lim_{s \to 0} sX(s)$$
(6.58)

$$\Rightarrow \lim_{t \to \infty} x(t) = \lim_{s \to 0} sX(s) \cdots \text{The final value theorem}$$
(6.59)

$$x(t) \longrightarrow h(t) \longrightarrow y(t)$$

$$X(s) H(s) Y(s)$$

$$Y(s) = H(s)X(s)$$

Figure 6.10 Block diagram of a system.

H(s): the system function or transfer function $s = j\omega$, $H(j\omega)$ is the frequency response of the LTI system. For a causal system, h(t) = 0 for t < 0 (Fig. 6.11). $\Rightarrow h(t)$ is a right-sided signal.

 \Rightarrow The ROC is the entire region in the *s*-plane to the right of the rightmost pole. Note:

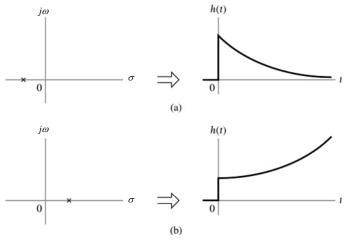
• Anticausal system h(t)

 \Rightarrow Its ROC is the region in the *s*-plane to the left of the leftmost pole.

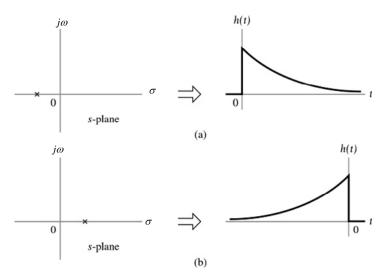
- An ROC to the right of the rightmost pole does not guarantee that the system is causal, only that the impulse response is right-sided.
- The Fourier transform of the impulse response for a stable LTI system exists.
 ⇒ For a stable system, the ROC of H(s) must include the jω-axis (Fig. 6.12).
- For a causal and stable LTI system with a rational system function, all poles must lie in the left half of the *s*-plane.

causal \rightarrow ROC is to the right of the rightmost pole.

stable \rightarrow ROC must include the $j\omega$ -axis.



■ Figure 6.11 The relationship between the locations of poles and the impulse response in a causal system. (a) A pole in the left half of the *s*-plane corresponds to an exponentially decaying impulse response. (b) A pole in the right half of the *s*-plane corresponds to an exponentially increasing impulse response. The system is unstable in this case.

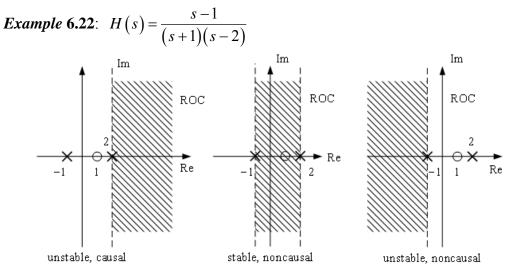


• Figure 6.12 The relationship between the locations of poles and the impulse response in a stable system. (a) A pole in the left half of the *s*-plane corresponds to a right-sided impulse response. (b) A pole in the right half of the *s*-plane corresponds to an left-sided impulse response. In this case, the system is noncausal.

Example 6.20:
$$h(t) = e^{-t}u(t) \Rightarrow H(s) = \frac{1}{s+1}$$
, $\operatorname{Re}\{s\} > -1$
 \Rightarrow causal and stable

Example 6.21:
$$H(s) = \frac{e^s}{s+1}$$
, $\operatorname{Re}\{s\} > -1$
 $e^{-t}u(t) \xleftarrow{\mathcal{L}} \frac{1}{s+1}$, $\operatorname{Re}\{s\} > -1$
 $e^{-(t+1)}u(t+1) \xleftarrow{\mathcal{L}} \frac{e^s}{s+1}$, $\operatorname{Re}\{s\} > -1$
 $\Rightarrow h(t) = e^{-(t+1)}u(t+1)$, zero for $t < -1$ but not for $t < 0$

 \Rightarrow not causal but is stable.



Example 6.23: Inverse Laplace transform with stability and causality constraints

$$H\left(s\right) = \frac{2}{s+3} + \frac{1}{s-2}$$

If the system is stable, then the pole at s = -3 contributes a right-sided term to the impulse response, while the pole at s = 2 contributes a left-sided term.

$$h(t) = 2e^{-3t}u(t) - e^{2t}u(-t)$$

If the system is causal, then both poles must contribute right-sided terms to the impulse response

$$h(t) = 2e^{-3t}u(t) + e^{2t}u(t)$$

1. System characterized by linear constant-coefficient differential equations

The system function has zeros at the solutions of

$$\sum_{k=0}^{M} b_k s^k = 0 \tag{6.62}$$

and poles at the solutions of

$$\sum_{k=0}^{N} a_k s^k = 0 \tag{6.63}$$

Note:

 With additional information such as stability or causality of the system, the ROC can be inferred and the corresponding impulse response can be obtained.

Example 6.24:

$$\frac{dy(t)}{dt} + 3y(t) = x(t)$$
$$\Rightarrow sY(s) + 3Y(s) = X(s) \Rightarrow H(s) = \frac{1}{s+3}$$

If the system is causal, the ROC is $\operatorname{Re}\{s\} > -3$, and the corresponding impulse response is

$$h(t) = e^{-3t}u(t)$$

If the system is noncausal, then the ROC is $\operatorname{Re}\{s\} < -3$, and the corresponding impulse response is

$$h(t) = -e^{-3t}u(-t)$$

- 2. System function for interconnections of LTI systems
 - (1) Parallel interconnection

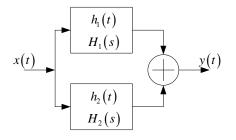


Figure 6.13 Parallel connection of two LTI systems.

$$h(t) = h_1(t) + h_2(t) \Longrightarrow H(s) = H_1(s) + H_2(s)$$

$$(6.64)$$

(2) Cascade interconnection

$$x(t) \qquad h_1(t) \qquad h_2(t) \qquad y(t) \\ H_1(s) \qquad H_2(s) \qquad h_2(s) \qquad h_2(t) \\ H_2(s) \qquad h_2(t) \qquad h_2(t) \\ H_2(t$$

Figure 6.14 Series connection of two LTI systems.

$$h(t) = h_1(t) * h_2(t) \Longrightarrow H(s) = H_1(s)H_2(s)$$
(6.65)

(3) Feedback interconnection

$$x(t) \xrightarrow{x_1(t)} \xrightarrow{h_1(t)} \xrightarrow{y_1(t)} \xrightarrow{y(t)} \xrightarrow{y(t)} \xrightarrow{y_2(t)} \xrightarrow{h_2(t)} \xrightarrow{x_2(t)} \xrightarrow{h_2(s)} \xrightarrow{x_2(t)} \xrightarrow{x_2(t)} \xrightarrow{h_2(s)} \xrightarrow{x_2(t)} \xrightarrow{x$$

Figure 6.15 Feedback interconnection of two LTI systems.

$$Y_{2}(s) = H_{2}(s)X_{2}(s) = H_{2}(s)Y_{1}(s) = H_{2}(s)Y(s)$$
(6.66)

$$Y(s) = H_1(s)X_1(s) = H_1(s)[X(s) - Y_2(s)]$$

= $H_1(s)X(s) - H_1(s)H_2(s)Y(s)$ (6.67)

$$\Rightarrow \frac{Y(s)}{X(s)} = H(s) = \frac{H_1(s)}{1 + H_1(s)H_2(s)}$$
(6.68)

3. Butterworth filters

$$|B(\omega)|^{2} = \frac{1}{1 + (\omega/\omega_{c})^{2N}} = \frac{1}{1 + (j\omega/j\omega_{c})^{2N}} = |B(j\omega)|^{2}$$
(6.69)

$$\left|B(j\omega)\right|^{2} = B(j\omega)B^{*}(j\omega)$$
(6.70)

Restricting the impulse response of the Butterworth filter to be real, we have

$$B^*(j\omega) = B(-j\omega) \tag{6.71}$$

$$B(j\omega)B(-j\omega) = 1/\left[1 + (j\omega/j\omega_c)^{2N}\right]$$
(6.72)

$$\therefore B(s)|_{s=j\omega} = B(j\omega) \tag{6.73}$$

$$\therefore B(s)B(-s) = 1/\left[1 + (s/j\omega_c)^{2N}\right]$$
(6.74)

The poles of B(s)B(-s) are the solutions of

$$1 + (s/j\omega_c)^{2N} = 0 (6.75)$$

$$\Rightarrow s_p = \left(-1\right)^{1/2N} \left(j\omega_c\right) \tag{6.76}$$

$$\Rightarrow \left| s_p \right| = \omega_c, \ \angle s_p = \frac{\pi (2k+1)}{2N} + \frac{\pi}{2}, \ k \text{ is an integer}$$
(6.77)

$$\Rightarrow s_p = \omega_c e^{j\left[\frac{\pi(2k+1)}{2N} + \frac{\pi}{2}\right]}$$
(6.78)

The positions of the poles of B(s)B(-s) for N = 1, 2, 3, and 6 are shown in Fig. 6.16.

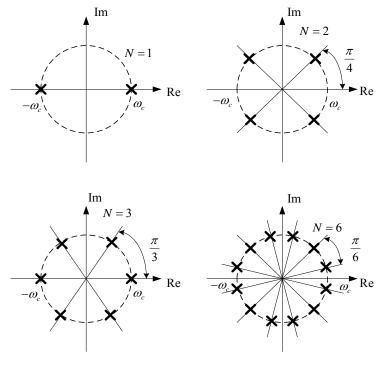


Figure 6.16 Position of B(s)B(-s) the poles of for N = 1, 2, 3, and 6.

- (1) The poles of B(s)B(-s) occurs in pairs, so that if there is a pole at $s = s_p$, then there is also a pole at $s = -s_p$.
- (2) To construct B(s), we choose one pole from each pair of poles.
- (3) If we restrict the system to be stable and causal, then the poles of

B(s) should be in the left-half plane.

(4)
$$B^2(s)\Big|_{s=0} = 1$$

$$N = 1: B(s) = \frac{\omega_c}{s + \omega_c}$$
(6.79)

$$N = 2: B(s) = \frac{\omega_c^2}{\left(s + \omega_c e^{j\frac{\pi}{4}}\right) \left(s + \omega_c e^{-j\frac{\pi}{4}}\right)} = \frac{\omega_c^2}{s^2 + \sqrt{2}\omega_c s + \omega_c^2}$$
(6.80)

$$N = 3: B(s) = \frac{\omega_c^3}{\left(s + \omega_c\right) \left(s + \omega_c e^{j\frac{\pi}{3}}\right) \left(s + \omega_c e^{-j\frac{\pi}{3}}\right)} = \frac{\omega_c^3}{s^3 + 2\omega_c s^2 + 2\omega_c^2 s + \omega_c^3}$$

(6.81)

The corresponding differential equations for the above three cases are:

$$N = 1: \quad \frac{dy(t)}{dt} + \omega_c y(t) = \omega_c x(t) \tag{6.82}$$

$$N = 2: \frac{d^2 y(t)}{dt^2} + \sqrt{2}\omega_c \frac{dy(t)}{dt} + \omega_c^2 y(t) = \omega_c^2 x(t)$$
(6.83)

$$N = 3: \ \frac{d^3 y(t)}{dt^3} + 2\omega_c \frac{d^2 y(t)}{dt^2} + 2\omega_c^2 \frac{dy(t)}{dt} + \omega_c^3 y(t) = \omega_c^3 x(t)$$
(6.84)

$$\frac{\omega_c}{s+\omega_c} = \frac{Y(s)}{X(s)} \Longrightarrow \omega_c X(s) = sY(s) + \omega_c Y(s)$$
(6.85)

$$\frac{\omega_c^2}{s^2 + \sqrt{2}\omega_c s + \omega_c^2} = \frac{Y(s)}{X(s)} \Longrightarrow \omega_c^2 X(s) = s^2 Y(s) + \sqrt{2}\omega_c s Y(s) + \omega_c^2 Y(s)$$

(6.86)

$$\frac{\omega_c^3}{s^3 + 2\omega_c s^2 + 2\omega_c^2 s + \omega_c^3} = \frac{Y(s)}{X(s)}$$

$$\Rightarrow \omega_c^3 X(s) = s^3 Y(s) + 2\omega_c s^2 Y(s) + 2\omega_c^2 s Y(s) + \omega_c^3 Y(s)$$
(6.87)

6-7 The Unilateral Laplace Transform

1. The unilateral Laplace transform of x(t) is defined as

$$\mathcal{X}(s) \triangleq \int_{0^{-}}^{\infty} x(t) e^{-st} dt \tag{6.88}$$

The bilateral Laplace transform

$$X(s) \triangleq \int_{-\infty}^{\infty} x(t) e^{-st} dt$$
(6.89)

when x(t) = 0 for t < 0, the unilateral and bilateral Laplace transforms are identical.

Note:

• The ROC for the unilateral Laplace transform is always a right-half plane (causal).

Example 6.25:

$$x(t) = \frac{t^{n-1}}{(n-1)!} e^{-at} u(t)$$
$$\mathcal{X}(s) = \frac{1}{(s+a)^n}, \operatorname{Re}\{s\} > -a$$

Example 6.26:

$$\begin{aligned} x(t) &= e^{-a(t+1)}u(t+1) \\ X(s) &= \frac{e^s}{s+a}, \ \operatorname{Re}\{s\} > -a \\ \mathcal{X}(s) &= \int_{0^-}^{\infty} e^{-a(t+1)}u(t+1)e^{-st}dt = \int_{0^-}^{\infty} e^{-a}e^{-(s+a)t}dt = e^{-a}\int_{\infty}^{\infty} \left\{ e^{-at}u(t) \right\}e^{-st}dt \\ &= \frac{e^{-a}}{s+a}, \ \operatorname{Re}\{s\} > -a \end{aligned}$$

The unilateral and bilateral Laplace transforms are distinctly different.

- 2. Most of the properties of the unilateral Laplace transform are the same as for the bilateral Laplace transform.
- 3. Differentiation property of the unilateral Laplace transform

$$\int_{0^{-}}^{\infty} \frac{dx(t)}{dt} e^{-st} dt = x(t) e^{-st} \Big|_{0^{-}}^{\infty} + s \int_{0^{-}}^{\infty} x(t) e^{-st} dt = s \mathcal{X}(s) - x(0^{-})$$
(Integration by parts)
(6.90)

 $\mathcal{X}(s)$ is the unilateral Laplace transform of x(t).

Similarly,

$$\int_{0^{-}}^{\infty} \frac{d^{2}x(t)}{dt^{2}} e^{-st} dt = s \int_{0^{-}}^{\infty} \frac{dx(t)}{dt} e^{-st} dt - x^{(1)} (0^{-})$$

= $s \left[s \mathcal{X}(s) - x(0^{-}) \right] - x^{(1)} (0^{-})$
= $s^{2} \mathcal{X}(s) - sx(0^{-}) - x^{(1)} (0^{-})$ (6.91)

The general form for the differentiation property is

$$\frac{d^{n}}{dt^{n}}x(t) \xleftarrow{\mathcal{L}_{u}}{s^{n}\mathcal{X}(s) - \frac{d^{n-1}}{dt^{n-1}}x(t)}\Big|_{t=0^{-}} - s\frac{d^{n-2}}{dt^{n-2}}x(t)\Big|_{t=0^{-}}$$

$$- \cdots - s^{n-2}\frac{d}{dt}x(t)\Big|_{t=0^{-}} - s^{n-1}x(0^{-})$$
(6.92)

where the subscript u in \mathcal{L}_u denotes the unilateral transform.

Example 6.27: $x(t) = e^{at}u(t)$

Apply the product rule for differentiation to obtain the derivative of x(t), $t > 0^-$:

$$\frac{d}{dt}x(t) = \frac{d}{dt}e^{at}u(t) = ae^{at}u(t) + \delta(t) \longleftrightarrow \frac{a}{s-a} + 1 = \frac{s}{s-a}$$

Using Eq. (6.90),

$$\frac{d}{dt}x(t) \longleftrightarrow s \frac{1}{s-a} + 0 = \frac{s}{s-a}$$

4. Integration property

$$\int_{-\infty}^{t} x(\tau) d\tau \xleftarrow{\mathcal{L}_{u}} \frac{\int_{-\infty}^{0^{-}} x(\tau) d\tau}{s} + \frac{\mathcal{X}(s)}{s}$$
(6.93)

Proof: Let $y(t) = \int_{-\infty}^{t} x(\tau) d\tau$. Then,

$$\frac{d}{dt}y(t) = x(t)$$

$$s\mathcal{Y}(s) - y(0^{-}) = \mathcal{X}(s)$$

$$\mathcal{Y}(s) = \frac{\mathcal{X}(s)}{s} + \frac{\int_{-\infty}^{0^{-}} x(\tau)d\tau}{s}$$

5. A primary use of the unilateral Laplace transform is in obtaining the solution of linear constant-coefficient differential equations with nonzero initial conditions.

Example 6.28:
$$\frac{d^2 y(t)}{dt^2} + 3 \frac{dy(t)}{dt} + 2y(t) = x(t), y(0^-) = 3, \frac{dy(0^-)}{dt} = -5$$

Let x(t) = 2u(t). Then we obtain

$$s^{2}\mathcal{Y}(s) - sy(0^{-}) - y^{(1)}(0^{-}) + 3s\mathcal{Y}(s) - 3y(0^{-}) + 2\mathcal{Y}(s) = \frac{2}{s}$$
$$\mathcal{Y}(s) = \frac{3s+4}{(s+1)(s+2)} + \frac{2}{s(s+1)(s+2)}$$

where $\mathcal{Y}(s)$ is the unilateral Laplace transform of y(t).

$$\Rightarrow \mathcal{Y}(s) = \frac{1}{s} - \frac{1}{s+1} + \frac{3}{s+2} \Rightarrow y(t) = \left[1 - e^{-t} + 3e^{-2t}\right]u(t)$$

6-8 Structures for Continuous-Time Filters

In Section 2-8 block diagram were employed to show the structure of continuous-time filter implementations as described by the corresponding differential (or integral) equations. The structure corresponding directly to the general differential equation in Eq. (6.60), with $a_N = 1$, was called the direct form and is shown in Fig. 6.17 for M = N, with 1/s denoting each integrator. Note that this structure consists effectively of the cascade of two subsystems.

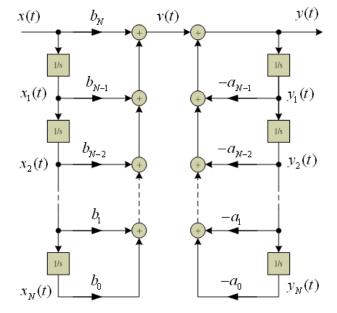


Figure 6.17 General continuous-time direct-form structure.

$$\frac{d^{N}v(t)}{dt^{N}} = \sum_{k=0}^{N} b_{k} \frac{d^{k}x(t)}{dt^{k}},$$
(6.94)

while the second subsystem realizes the differential equation

$$\sum_{k=0}^{N} a_k \frac{d^k y(t)}{dt^k} = \frac{d^N v(t)}{dt^N},$$
(6.95)

with $a_N = 1$. Calling these subsystems $H_1(s)$ and $H_2(s)$, respectively, we thus have

$$H(s) = H_1(s)H_2(s) = B(s)/A(s), \qquad (6.96)$$

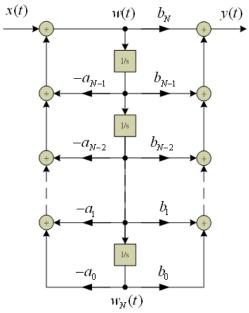
where

$$H_1(s) = \frac{B(s)}{s^N} = \frac{\sum_{k=0}^N b_k s^k}{s^N}$$
(6.97)

and

$$H_{2}(s) = \frac{s^{N}}{A(s)} = \frac{s^{N}}{\sum_{k=0}^{N} a_{k} s^{k}}$$
(6.98)

As noted in Section 2-8, the direct form in Fig. 6.17 is not canonical because the number of integrators (2*N*) is not minimum. Reversing the order of $H_1(s)$ and $H_2(s)$ and eliminating the *N* redundant integrators, we produce the canonical direct-form-II, as shown in Fig. 6.18 Note that, in addition to *N* integrators, this canonical form includes 2*N*+1 multipliers (amplifiers), in general, for an *N*th-order filter with M = N.



■ Figure 6.18 General continuous-time direct-form-II.

Direct-form-II implementations of the first-order LPF and HPF from

$$H(s) = \frac{1}{1 + s / \omega_b} = \frac{\omega_b}{s + \omega_b}, \qquad (6.99)$$

and

$$H(s) = \frac{s}{s + \omega_b},\tag{6.100}$$

are shown in Fig. 6.19 Note that the signs of the feedback multipliers $(-\omega_b)$ and the corresponding terms in the denominators of H(s) are different, as opposed to the feedforward coefficients $(\omega_b \text{ and } 1)$ and the corresponding numerator terms.

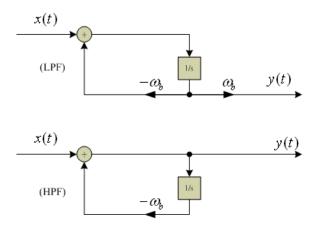


Figure 6.19 First-order LPF and HPF direct-form-II structures.

There are many other structures (canonical and otherwise) that are useful for implementing continuous-time (analog) filters. These structures have various desirable properties such as modularity, reduced sensitivity to component variation, and/or suitability for integrated-circuit realization. Although we will consider only active filter implementations in this section, passive-circuit realizations are possible. Two basic modulator structures are the parallel form and cascade form. To derive the parallel form, we expand H(s) in the partial-fraction expansion (assuming no multiple poles)

$$H(s) = b_N + \sum_{k=1}^{N} \frac{r_k}{s + s_k},$$
(6.101)

where $b_N = 0$ if M < N. This form for H(s) implies a parallel combination of N first-order filters. In general, the poles $-s_k$ and reduces r_k are complex-valued, complex multipliers would be required in the corresponding implementation. In particular, assuming that h(t) is real-valued, H(s) can be written as

$$H(s) = b_N + \sum_{k=1}^{L} \left(\frac{r_k}{s+s_k} + \frac{r_k^*}{s+s_k^*} \right) + \sum_{k=2L+1}^{N} \frac{r_k}{s+s_k},$$
(6.102)

where s_k , k=1,...,L ($L \le N/2$) are complex-valued and s_k , k=2L+1,...,N, are real-valued. To avoid the unnecessary complication of complex multipliers, we combine the terms in the first summation to obtain

$$H(s) = b_N + \sum_{k=1}^{L} \frac{\gamma_{1k} s + \gamma_{0k}}{s^2 + \alpha_{1k} s + \alpha_{0k}} + \sum_{k=2L+1}^{N} \frac{r_k}{s + s_k},$$
(6.103)

where

$$\alpha_{1k} = 2 \operatorname{Re}\{s_k\} \qquad \alpha_{0k} = |s_k|^2$$

$$\gamma_{1k} = 2 \operatorname{Re}\{r_k\} \qquad \gamma_{0k} = 2 \operatorname{Re}\{s_k^* r_k\}$$

Hence all of the coefficients in Eq. (6.103) are real-valued. Using direct-form-II networks to realize each of the terms in Eq. (6.103), we produce parallel form II, which is shown in Fig. 6.20 for N odd and L=(N-1)/2. Note that the parallel form is also canonical since it has N integrators and 2N-1 multipliers, in general, for M = N. If several poles are real-valued (that is, if $N-2L \ge 2$), some or all of the associated first-order terms in Eq. (6.103) are often combined into second-order terms to produce additional second-order sections in the parallel form, which increases the modularity of the corresponding circuit realization.

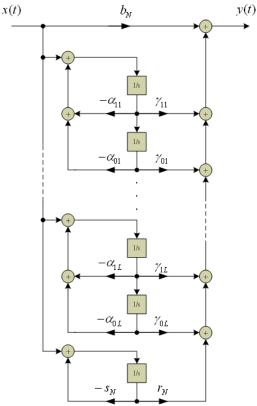


Figure 6.20 *N*th-order parallel-form-II structure for *N* odd.

To obtain the cascade form, we instead factor the numerator and denominator of H(s) into products of first-order terms of the form

$$H(s) = b_M \frac{\prod_{k=1}^{M} (s + v_k)}{\prod_{k=1}^{N} (s + s_k)}.$$
 (6.104)

Assuming that $a_N = 1$, as before. (Multiple poles and/or zeros are allowed.) This form for H(s) implies a cascade combination of first-order subfilters, but again complex-valued poles $-s_k$ and/or zeros $-v_k$ would necessitate complex multipliers. Therefore we rewrite H(s) as

$$H(s) = b_M \frac{\prod_{k=1}^{K} (s + v_k)(s + v_k^*) \prod_{k=2K+1}^{M} (s + v_k)}{\prod_{k=1}^{L} (s + s_k)(s + s_k^*) \prod_{k=2L+1}^{N} (s + s_k)},$$
(6.105)

where v_k , k=1,...,K ($K \le M/2$), and s_k , k=1,...,L ($L \le N/2$), are complex-valued and v_k , k = 2K+1,...,M, and s_k , k = 2L+1,...,N, are real-valued. The complex factors are then combined to produce

$$H(s) = b_{M} \frac{\prod_{k=1}^{K} (s^{2} + \beta_{1k}s + \beta_{0k}) \prod_{k=2K+1}^{M} (s + v_{k})}{\prod_{k=1}^{L} (s^{2} + \alpha_{1k}s + \alpha_{0k}) \prod_{k=2L+1}^{N} (s + s_{k})},$$
(6.106)

where

$$\alpha_{1k} = -2 \operatorname{Re}\{s_k\} \quad \alpha_{0k} = |s_k|^2 \beta_{1k} = -2 \operatorname{Re}\{r_k\} \quad \beta_{0k} = |v_k|^2$$

Therefore, since all of the coefficients in this expression are real-valued, H(s) can be implemented as a cascade of first- and second-order sections with real multipliers. Again, if there are several real-valued poles and/or zeros (i.e., if $N-2L\geq 2$ and/or $M-2K\geq 2$), the corresponding factors are usually combined in pairs to produce additional second-order sections. Realizing the resulting first- and second-order sections using direct-form-II networks, we produce cascade form II, shown in Fig. 6.21 for M = N (odd) and L=(N-1)/2. Note that the cascade form is also, in general, canonical.

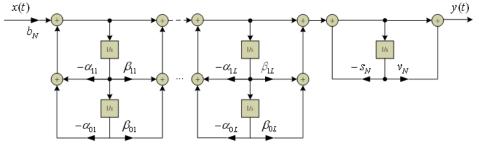


Figure 6.21 *N*th-order cascade-form-II structure for *N* odd.

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